

Low energy neutrino experiments sensitivity to physics beyond the Standard Model

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with

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based on hep-ph/0508299 and hep-ph/0702175

Outline

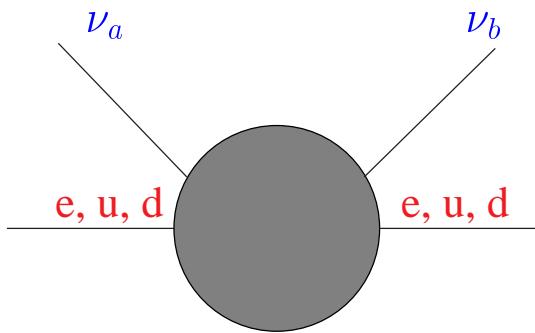
- Motivations
- Non-standard contributions to ν - N coherent scattering
- Experimental sensitivity
- Sensitivity to specific NSI scenarios:
 Z' , leptoquark and R-parity breaking SUSY
- Summary

Non Standard Interactions (NSI)

Most extensions of the SM, in particular neutrino mass theories, predict neutral current non-standard interactions (NSI) of neutrinos which can be either flavor preserving (NU – non-universal) or flavor-changing (FC).

NSI effective Lagragian form:

$$\mathcal{L}_{eff}^{NSI} = - \sum_{\alpha\beta fP} \varepsilon_{\alpha\beta}^{fP} 2\sqrt{2} G_F (\bar{\nu}_\alpha \gamma_\rho L \nu_\beta) (\bar{f} \gamma^\rho P f)$$



Here $\alpha, \beta = e, \mu, \tau$; $f = e, u, d$; $P = L, R$; $L = (1 - \gamma_5)/2$; $R = (1 + \gamma_5)/2$

Non Standard Interactions (NSI)

Non-standard neutral current neutrino interactions may arise:

- from a non trivial non-unitary lepton mixing matrix
Schechter & Valle'80
- in models where neutrino masses are "calculable" from radiative corrections
Zee'80, Babu'88
- in SUSY models with broken R-parity
see review by Hirsch & Valle [hep-ph/0405015] and refs therein
- in unified SUSY models as a renormalization effect
Hall, Kostelecky & Raby'86
- ... some other models, like left-right models, etc ...

Predictions:

In most models NSI contributions are expected to be small, e.g. being suppressed by the smallness of neutrino masses, however in some models NSI is not strongly restricted or suppressed

Example: \mathbb{R}_p parity violating SUSY

Non-standard neutrino-electron and neutrino quark interactions:

$$\mathcal{L} = \lambda_{ijk} \tilde{e}_R^{k*} (\bar{\nu}_L^i)^c e_L^j + \lambda'_{ijk} \tilde{d}_L^j \bar{d}_R^k \nu_L^i + \dots$$

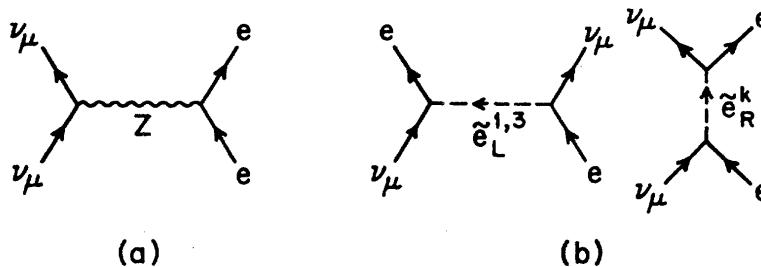
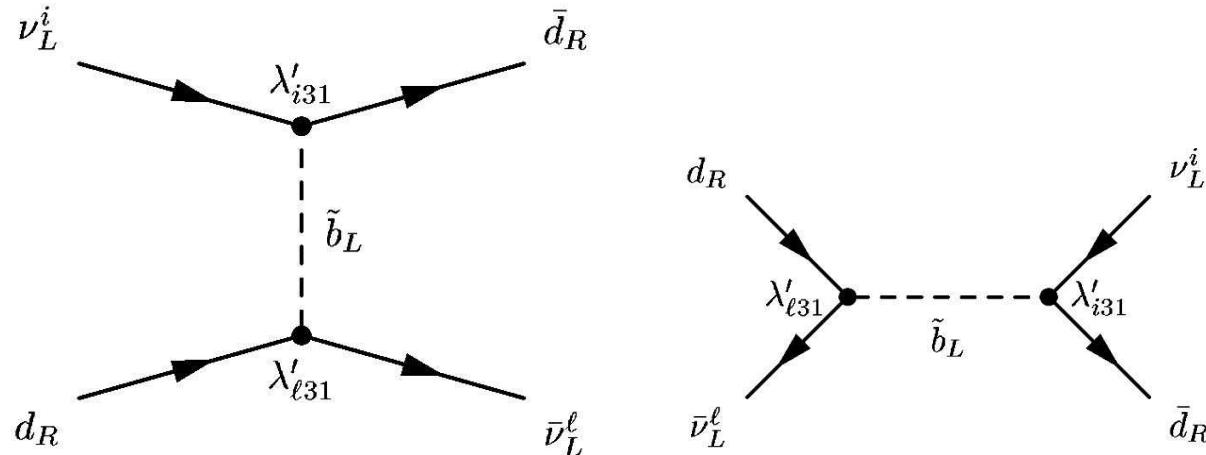


FIG. 2. Feynman diagrams for $\nu_\mu e$ scattering from (a) the standard model, and (b) the R -breaking interactions.

Barger, Giudice & Han'89



See e.g. Roulet'91, Amanik et al'05

Applications

- neutrino oscillations in matter: solar, atmospheric, supernovae
..., Guzzo et al'91, Fornengo et al'02, Friedland et al'04'05, Miranda et al'04
- supernovae explosion
Freedman et al'77, Fuller et al'87'88, Amanik et al'04'06, Esteban-Pretel et al'07
- ν scattering experiments
..., Barger et al'91, Davidson et al'03
- LEP (ILC)
Berezhiani & Rossi'02, Davidson et al'03
- ...

Solar ν oscillations and NSI

$$H_{\text{NSI}} = \sqrt{2}G_F N_f \begin{pmatrix} 0 & \varepsilon \\ \varepsilon & \varepsilon' \end{pmatrix}.$$

with,

$$\varepsilon = -\sin \theta_{23} \varepsilon_{e\tau}^{fV} \quad \varepsilon' = \sin^2 \theta_{23} \varepsilon_{\tau\tau}^{fV} - \varepsilon_{ee}^{fV}$$

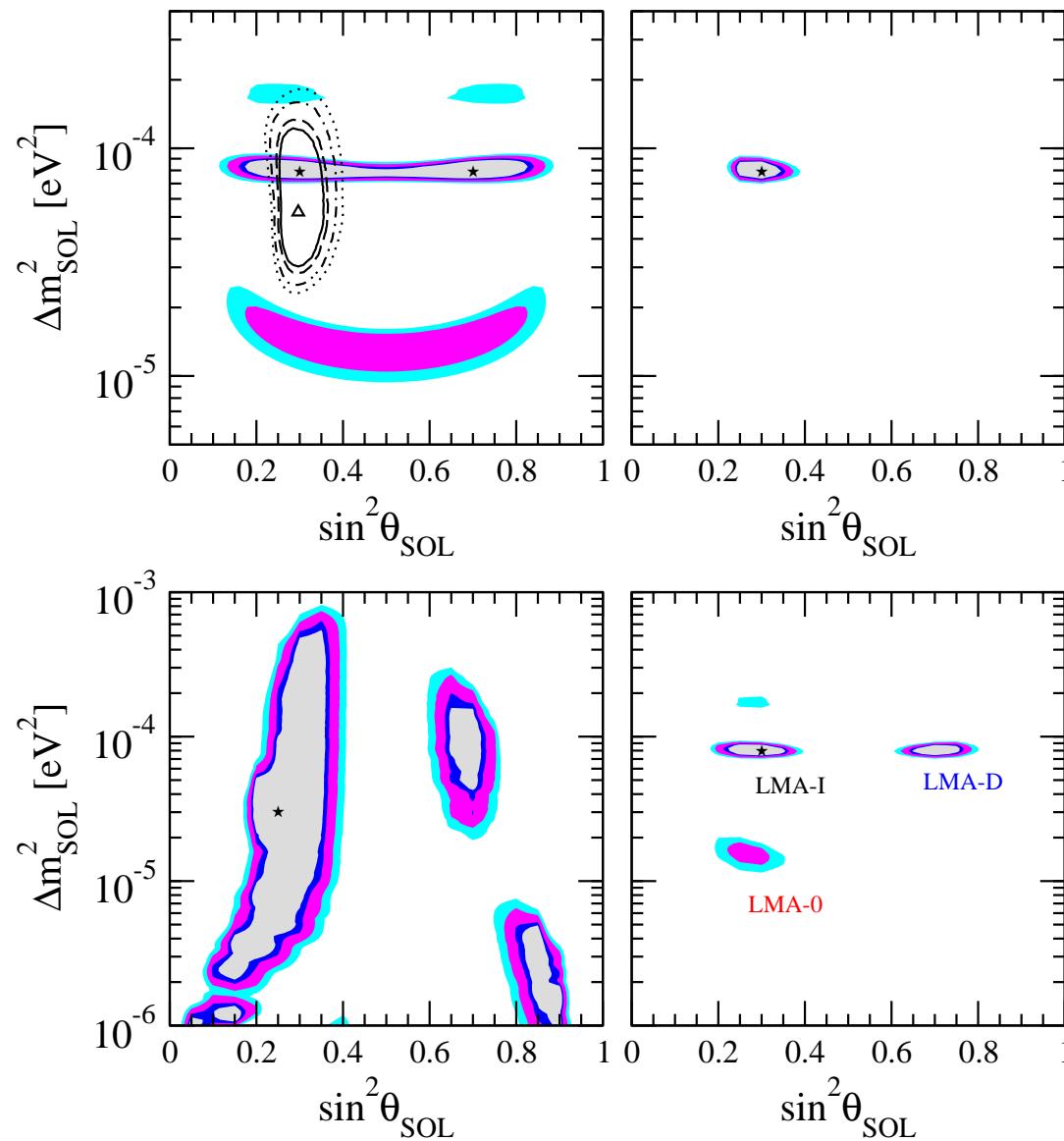
and

$$\varepsilon_{\tau\tau}^{fV} = \varepsilon_{\tau\tau}^{fL} + \varepsilon_{\tau\tau}^{fR}$$

Note:

ν oscillations are sensitive mainly to vector couplings because matter is not polarized.

Solar + KamLAND without and with NSI



Current bounds on NSI couplings

Bounds on NSI couplings come from

- ν -scattering experiments: LSND, CHARM, NuTeV, MUNU, MINOS
Barger et al'91, Davidson et al'03, Barranco et al'05
Friedland et al'06
- $e^- e^+ \rightarrow \nu \bar{\nu} \gamma$ measured at LEP
Berezhiani & Rossi'02
- analysis of atmospheric neutrino data
Fornengo et al'02, Friedland et al'04'05
- lepton flavor violating interactions, appeared at loop level from NSI,
like μ capture by nuclei
Davidson et al'03
- Invisible Z-boson decay width including loop corrections due to NSI
Davidson et al'03

Current bounds on NU NSI ν - q couplings from Davidson et al'03

| vertex | current limits | future limit |
|---|---|--|
| $(\bar{u}\gamma^\rho P u)(\bar{\nu}_\tau \gamma_\rho L \nu_\tau)$ | $ \varepsilon_{\tau\tau}^{uL} < 1.4$ $ \varepsilon_{\tau\tau}^{uR} < 3$ $(\Gamma_{inv})^*)$ | $-0.3 < \varepsilon_{\tau\tau}^{uL} < 0.25$ $-0.25 < \varepsilon_{\tau\tau}^{uR} < 0.3$ KamLAND and SNO/SK |
| $(\bar{d}\gamma^\rho L d)(\bar{\nu}_\tau \gamma_\rho L \nu_\tau)$ | $ \varepsilon_{\tau\tau}^{dL} < 1.1$ $ \varepsilon_{\tau\tau}^{dR} < 6$ $(\Gamma_{inv})^*)$ | $-0.25 < \varepsilon_{\tau\tau}^{dL} < 0.3$ $-0.3 < \varepsilon_{\tau\tau}^{dR} < 0.25$ KamLAND and SNO/SK |
| $(\bar{u}\gamma^\rho P u)(\bar{\nu}_\mu \gamma_\rho L \nu_\mu)$ | $ \varepsilon_{\mu\mu}^{uL} < 0.003$ $-0.008 < \varepsilon_{\mu\mu}^{uR} < 0.003$ NuTeV | $ \varepsilon_{\mu\mu}^{uL} < 0.001$ $ \varepsilon_{\mu\mu}^{uR} < 0.002$ s_W^2 in DIS at ν Factory |
| $(\bar{d}\gamma^\rho P d)(\bar{\nu}_\mu \gamma_\rho L \nu_\mu)$ | $ \varepsilon_{\mu\mu}^{dL} < 0.003$ $-0.008 < \varepsilon_{\mu\mu}^{dR} < 0.015$ NuTeV | $ \varepsilon_{\mu\mu}^{dL} < 0.0009$ $ \varepsilon_{\mu\mu}^{dR} < 0.005$ s_W^2 in DIS at ν Factory |
| $(\bar{u}\gamma^\rho P u)(\bar{\nu}_e \gamma_\rho L \nu_e)$ | $-1 < \varepsilon_{ee}^{uL} < 0.3$ $-0.4 < \varepsilon_{ee}^{uR} < 0.7$ CHARM | $ \varepsilon_{ee}^{uL} < 0.001$ $ \varepsilon_{ee}^{uR} < 0.002$ s_W^2 in DIS at ν Factory |
| $(\bar{d}\gamma^\rho P d)(\bar{\nu}_e \gamma_\rho L \nu_e)$ | $-0.3 < \varepsilon_{ee}^{dL} < 0.3$ $-0.6 < \varepsilon_{ee}^{dR} < 0.5$ CHARM | $ \varepsilon_{ee}^{dL} < 0.0009$ $ \varepsilon_{ee}^{dR} < 0.005$ s_W^2 in DIS at ν Factory |

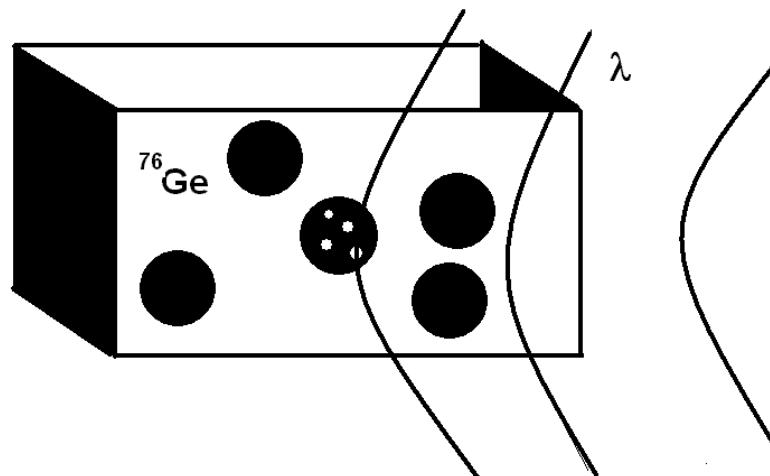
Current bounds on FC NSI ν - q couplings from Davidson et al'03

| vertex | current limits | future limit |
|--|--|--|
| $(\bar{u}\gamma^\rho P u)(\bar{\nu}_\tau \gamma_\rho L \nu_\mu)$ | $ \varepsilon_{\tau\mu}^{uP} < 0.05$ NuTeV | $ \varepsilon_{\tau\mu}^{uP} < 0.03$ s_W^2 in DIS at ν Factory |
| $(\bar{d}\gamma^\rho P d)(\bar{\nu}_\tau \gamma_\rho L \nu_\mu)$ | $ \varepsilon_{\tau\mu}^{dP} < 0.05$ NuTeV | $ \varepsilon_{\tau\mu}^{dP} < 0.03$ s_W^2 in DIS at ν Factory |
| $(\bar{u}\gamma^\rho P u)(\bar{\nu}_\mu \gamma_\rho L \nu_e)$ | $ \varepsilon_{\mu e}^{uP} < 7.7 \times 10^{-4}$ (Ti $\mu \rightarrow$ Ti e) *) | |
| $(\bar{d}\gamma^\rho P d)(\bar{\nu}_\mu \gamma_\rho L \nu_e)$ | $ \varepsilon_{\mu e}^{dP} < 7.7 \times 10^{-4}$ (Ti $\mu \rightarrow$ Ti e) *) | |
| $(\bar{u}\gamma^\rho P u)(\bar{\nu}_\tau \gamma_\rho L \nu_e)$ | $ \varepsilon_{\tau e}^{uP} < 0.5$ CHARM | $ \varepsilon_{\tau e}^{uP} < 0.03$ s_W^2 in DIS at ν Factory |
| $(\bar{d}\gamma^\rho P d)(\bar{\nu}_\tau \gamma_\rho L \nu_e)$ | $ \varepsilon_{\tau e}^{dP} < 0.5$ CHARM | $ \varepsilon_{\tau e}^{dP} < 0.03$ s_W^2 in DIS at ν Factory |

- All these bounds are derived taking one parameter at a time!
- NSI couplings with ν_μ are already strongly restricted
- See Davidson et al'03 [hep-ph/0302093] for NSI neutrino-electron couplings bounds

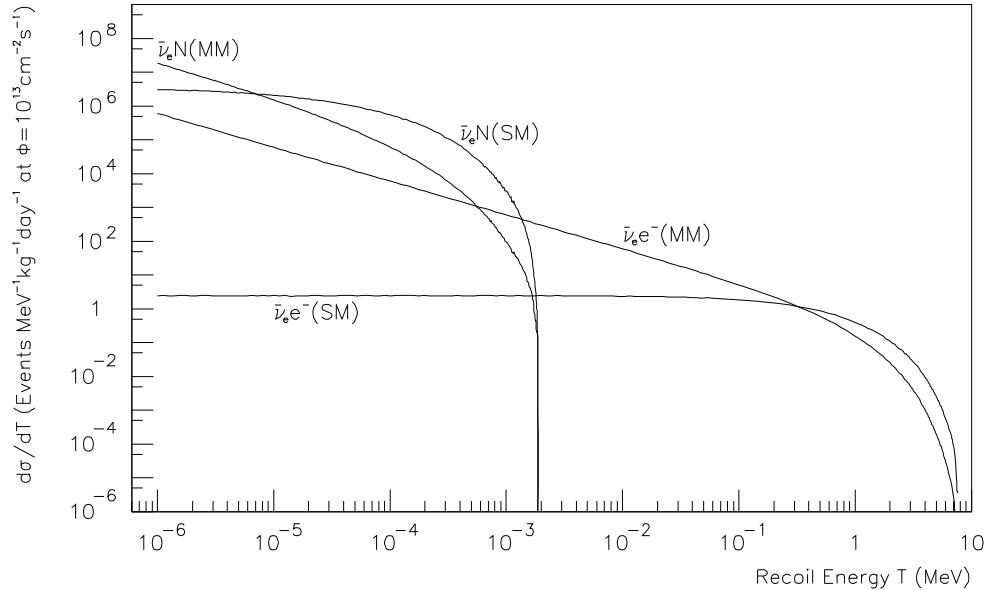
Coherent neutrino scattering off nuclei

- Good statistics due to quadratic coherent enhancement
- Sensitivity to ν -quark couplings
- Coherent scattering if the momentum transfer, Q , is small, $QR < 1$ (R is radius of nucleus): $\Rightarrow \nu$ -s doesn't "see" structure of nucleus!
- For most of nuclei: $1/R \sim 25 - 150$ MeV
- Well satisfied for most neutrino sources like supernovae, solar, reactor and artificial neutrino sources
- Planned experiments to measure coherent ν - N scattering: NOSTOS, TEXONO ... and many proposals
- Experimentally difficult: very low energy threshold



Proposed experiments to measure coherent ν - N scattering

- **TEXONO**: 1kg of germanium, reactor neutrinos hep-ex/0511001
 - NOSTOS: spherical TPC detector, 10 ton of Xenon astro-ph/0511470
 - Stopped-pion neutrino beam and kg-to-ton mass detector
hep-ex/0511042
 - more ideas in the past,
 - superconducting detector (Drukier & Stodolsky'84)
 - acoustic (Krauss'91)
 - cryogenic (Oberauer'02)



Neutrino-nuclei interaction

$$\mathcal{L}_{\nu Hadron}^{NSI} = -\frac{G_F}{\sqrt{2}} \sum_{\substack{q=u,d \\ \alpha,\beta=e,\mu,\tau}} [\bar{\nu}_\alpha \gamma^\mu (1-\gamma^5) \nu_\beta] \left(\varepsilon_{\alpha\beta}^{qL} [\bar{q} \gamma_\mu (1-\gamma^5) q] + \varepsilon_{\alpha\beta}^{qR} [\bar{q} \gamma_\mu (1+\gamma^5) q] \right),$$

$$\mathcal{L}_{\nu Hadron}^{NC} = -\frac{G_F}{\sqrt{2}} \sum_{\substack{q=u,d \\ \alpha,\beta=e,\mu,\tau}} [\bar{\nu}_\alpha \gamma^\mu (1-\gamma^5) \nu_\beta] \left(f_{\alpha\beta}^{qL} [\bar{q} \gamma_\mu (1-\gamma^5) q] + f_{\alpha\beta}^{qR} [\bar{q} \gamma_\mu (1+\gamma^5) q] \right),$$

$$f_{\alpha\alpha}^{uL} = \rho_{\nu N}^{NC} \left(\frac{1}{2} - \frac{2}{3} \hat{\kappa}_{\nu N} \hat{s}_Z^2 \right) + \lambda^{uL} + \varepsilon_{\alpha\alpha}^{uL}$$

$$f_{\alpha\alpha}^{dL} = \rho_{\nu N}^{NC} \left(-\frac{1}{2} + \frac{1}{3} \hat{\kappa}_{\nu N} \hat{s}_Z^2 \right) + \lambda^{dL} + \varepsilon_{\alpha\alpha}^{dL}$$

$$f_{\alpha\alpha}^{uR} = \rho_{\nu N}^{NC} \left(-\frac{2}{3} \hat{\kappa}_{\nu N} \hat{s}_Z^2 \right) + \lambda^{uR} + \varepsilon_{\alpha\alpha}^{uR}$$

$$f_{\alpha\alpha}^{dR} = \rho_{\nu N}^{NC} \left(\frac{1}{3} \hat{\kappa}_{\nu N} \hat{s}_Z^2 \right) + \lambda^{dR} + \varepsilon_{\alpha\alpha}^{dR}$$

ν -N coherent scattering

$$\frac{d\sigma}{dT} = \frac{G_F^2 M}{2\pi} \left\{ (G_V + G_A)^2 + (G_V - G_A)^2 \left(1 - \frac{T}{E_\nu}\right)^2 - (G_V^2 - G_A^2) \frac{MT}{E_\nu^2} \right\}$$

$$\begin{aligned} G_V &= \left[\left(g_V^p + 2\varepsilon_{ee}^{uV} + \varepsilon_{ee}^{dV} \right) Z + \left(g_V^n + \varepsilon_{ee}^{uV} + 2\varepsilon_{ee}^{dV} \right) N \right] F_{nucl}^V(Q^2) \\ G_A &= \left[\left(g_A^p + 2\varepsilon_{ee}^{uA} + \varepsilon_{ee}^{dA} \right) (Z_+ - Z_-) + \left(g_A^n + \varepsilon_{ee}^{uA} + 2\varepsilon_{ee}^{dA} \right) (N_+ - N_-) \right] F_{nucl}^A(Q^2) \end{aligned}$$

$$\begin{aligned} \frac{d\sigma}{dT}(E_\nu, T) &= \frac{G_F^2 M}{\pi} \left(1 - \frac{MT}{2E_\nu^2}\right) \times \\ &\times \left\{ \left[Z(g_V^p + 2\varepsilon_{ee}^{uV} + \varepsilon_{ee}^{dV}) + N(g_V^n + \varepsilon_{ee}^{uV} + 2\varepsilon_{ee}^{dV}) \right]^2 + \right. \\ &\left. + \sum_{\alpha=\mu,\tau} \left[Z(2\varepsilon_{\alpha e}^{uV} + \varepsilon_{\alpha e}^{dV}) + N(\varepsilon_{\alpha e}^{uV} + 2\varepsilon_{\alpha e}^{dV}) \right]^2 \right\} \end{aligned}$$

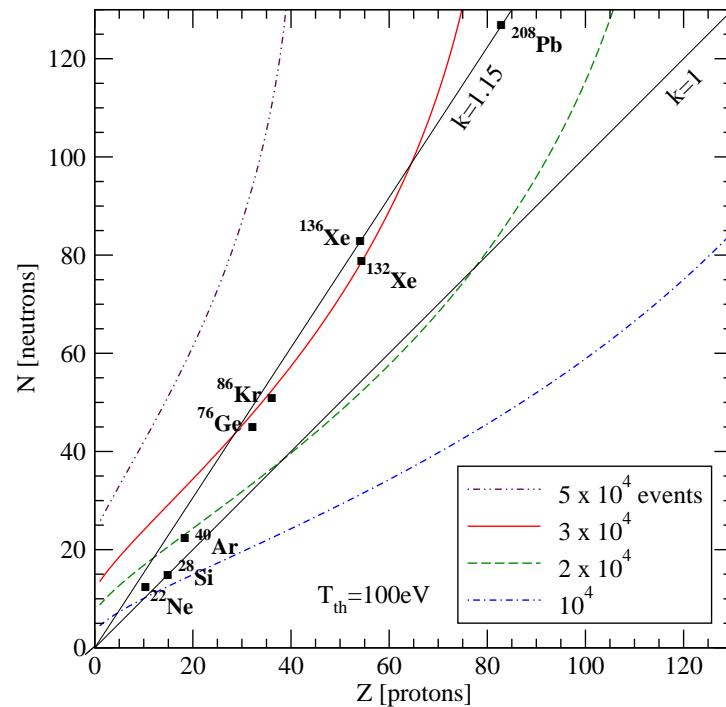
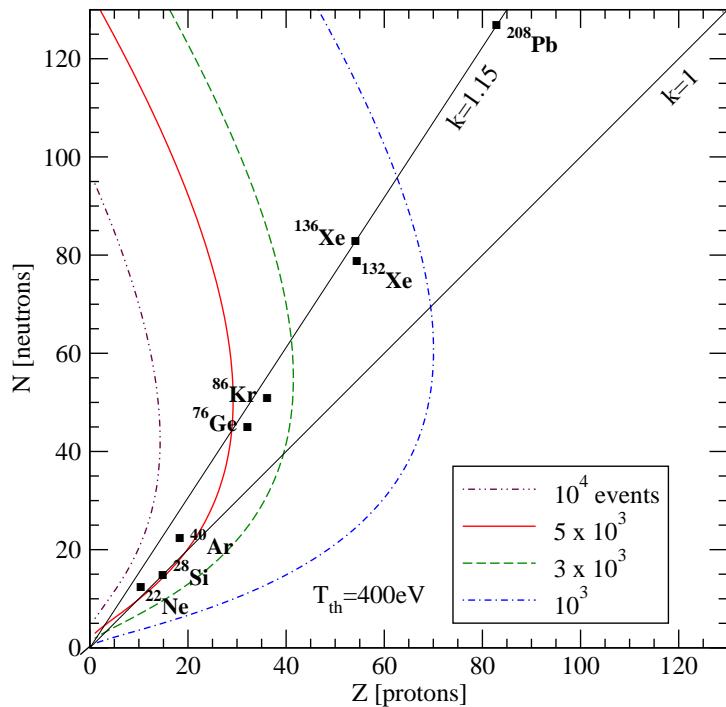
- Axial couplings contribution is zero or can be neglected
- Coherent enhancement of cross section
- Degeneracy in determination of NSI parameters

Resolving degeneracy

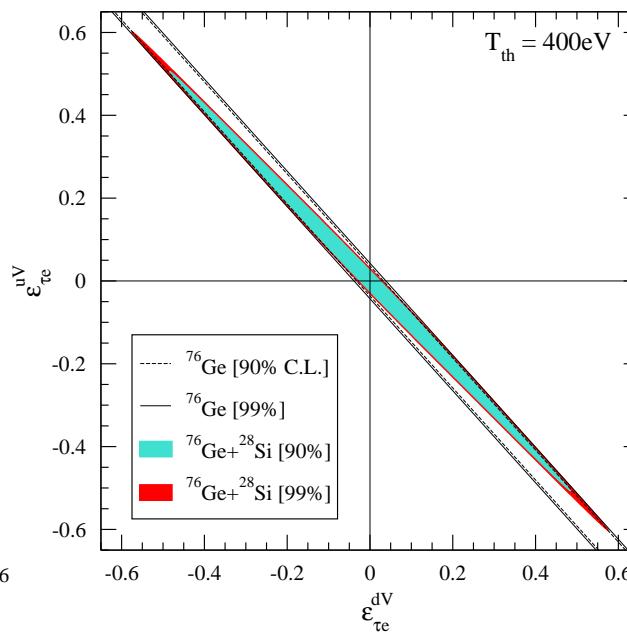
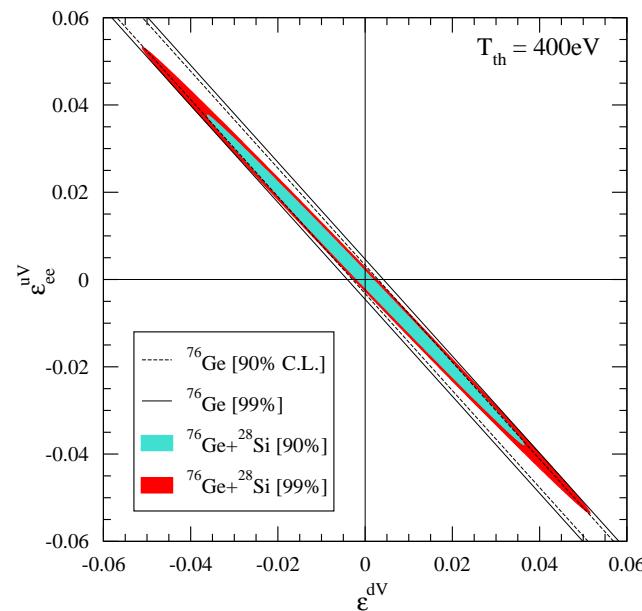
$$\left[Z(g_V^p + 2\varepsilon_{ee}^{uV} + \varepsilon_{ee}^{dV}) + N(g_V^n + \varepsilon_{ee}^{uV} + 2\varepsilon_{ee}^{dV}) \right]^2 = [Zg_V^p + Ng_V^n]^2$$

$$\varepsilon_{ee}^{uV}(A+Z) + \varepsilon_{ee}^{dV}(A+N) = \text{const.}$$

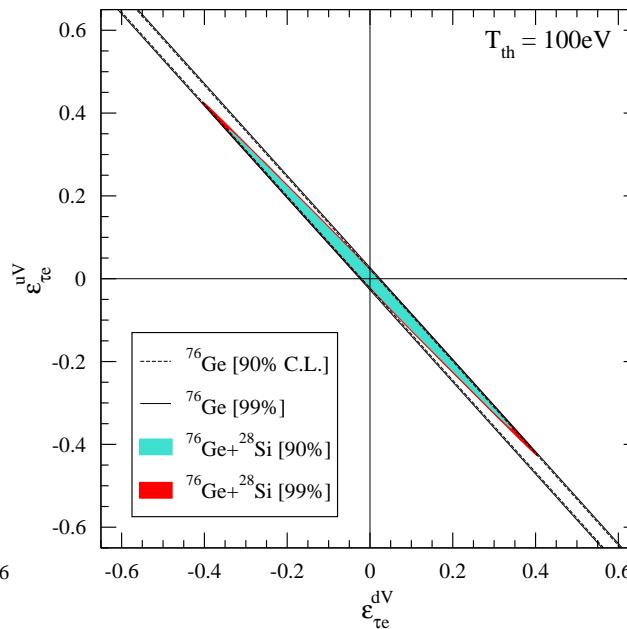
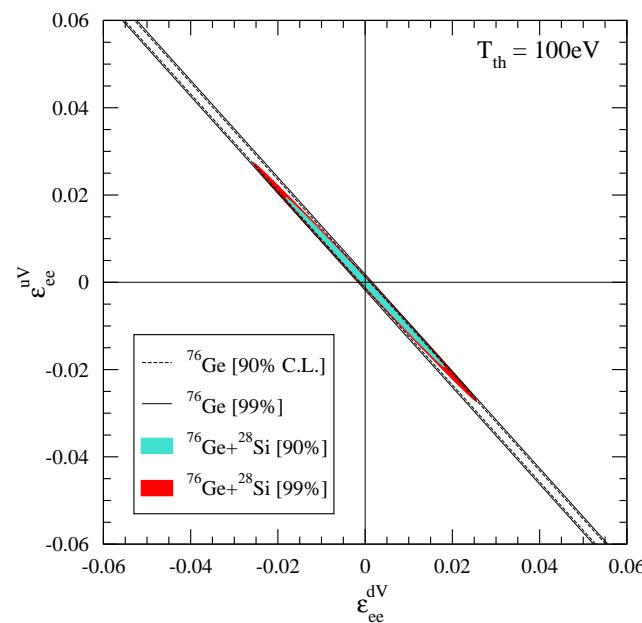
Solution: take two targets with **maximally different** $k = (A+N)/(A+Z)$



Estimated bounds on NSI from TEXONO-like experiment (Ge+Si)



| ${}^{76}\text{Ge} + {}^{28}\text{Si}$ $T_{th}=400\text{eV}$ |
|---|
| $ \epsilon_{ee}^{dV} < 0.036$ |
| $ \epsilon_{ee}^{uV} < 0.038$ |
| $ \epsilon_{\tau e}^{dV} < 0.48$ |
| $ \epsilon_{\tau e}^{uV} < 0.50$ |



| ${}^{76}\text{Ge} + {}^{28}\text{Si}$ $T_{th}=100\text{eV}$ |
|---|
| $ \epsilon_{ee}^{dV} < 0.018$ |
| $ \epsilon_{ee}^{uV} < 0.019$ |
| $ \epsilon_{\tau e}^{dV} < 0.34$ |
| $ \epsilon_{\tau e}^{uV} < 0.37$ |

Present and future bounds on NSI

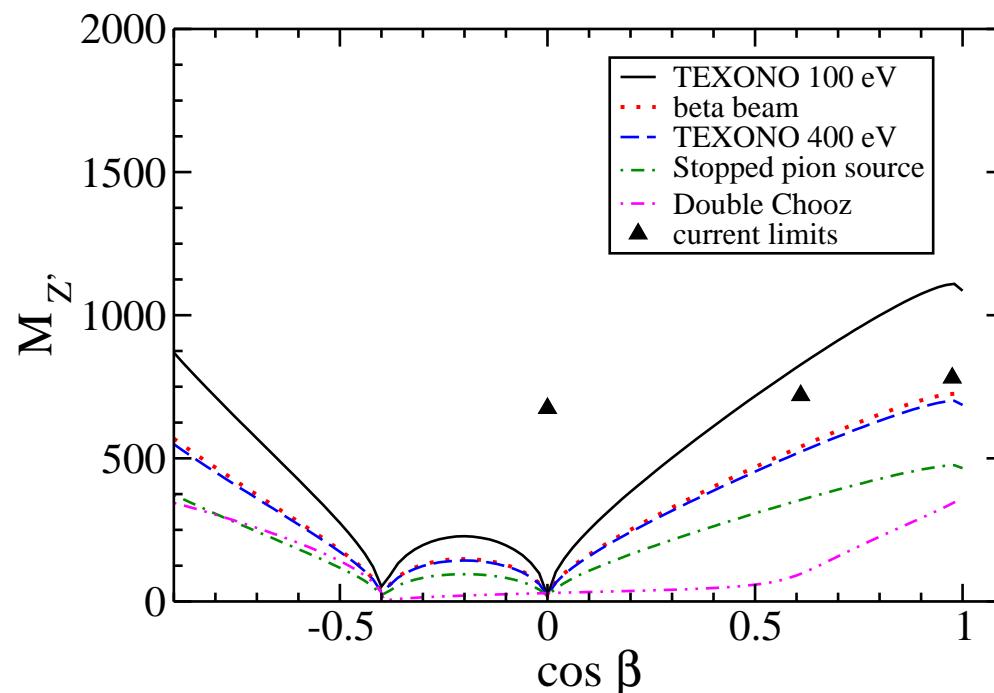
One parameter analysis to compare coherent scattering sensitivity with present bounds and ν Factory sensitivity (taken from Davidson et al'03)

| | Present Limits | ν Factory | ${}^{76}\text{Ge}$ $T_{th}=400\text{eV}$ (${}^{76}\text{Ge}$ $T_{th}=100\text{eV}$) | ${}^{76}\text{Ge}+{}^{28}\text{Si}$ $T_{th}=400\text{eV}$ (${}^{76}\text{Ge}+{}^{28}\text{Si}$ $T_{th}=100\text{eV}$) |
|--------------------------|------------------------------------|-----------------------------------|--|--|
| ϵ_{ee}^{dV} | $-0.5 < \epsilon_{ee}^{dV} < 1.2$ | $ \epsilon_{ee}^{dV} < 0.002$ | $ \epsilon_{ee}^{dV} < 0.003$ $(\epsilon_{ee}^{dV} < 0.001)$ | $ \epsilon_{ee}^{dV} < 0.002$ $(\epsilon_{ee}^{dV} < 0.001)$ |
| $\epsilon_{\tau e}^{dV}$ | $ \epsilon_{\tau e}^{dV} < 0.78$ | $ \epsilon_{\tau e}^{dV} < 0.06$ | $ \epsilon_{\tau e}^{dV} < 0.032$ $(\epsilon_{\tau e}^{dV} < 0.020)$ | $ \epsilon_{\tau e}^{dV} < 0.024$ $(\epsilon_{\tau e}^{dV} < 0.017)$ |
| ϵ_{ee}^{uV} | $-1.0 < \epsilon_{ee}^{uV} < 0.61$ | $ \epsilon_{ee}^{uV} < 0.002$ | $ \epsilon_{ee}^{uV} < 0.003$ $(\epsilon_{ee}^{uV} < 0.001)$ | $ \epsilon_{ee}^{uV} < 0.002$ $(\epsilon_{ee}^{uV} < 0.001)$ |
| $\epsilon_{\tau e}^{uV}$ | $ \epsilon_{\tau e}^{uV} < 0.78$ | $ \epsilon_{\tau e}^{uV} < 0.06$ | $ \epsilon_{\tau e}^{uV} < 0.036$ $(\epsilon_{\tau e}^{uV} < 0.023)$ | $ \epsilon_{\tau e}^{uV} < 0.023$ $(\epsilon_{\tau e}^{uV} < 0.018)$ |

Extra heavy neutral gauge boson Z'

See Gonzalez-Garcia & Valle'90

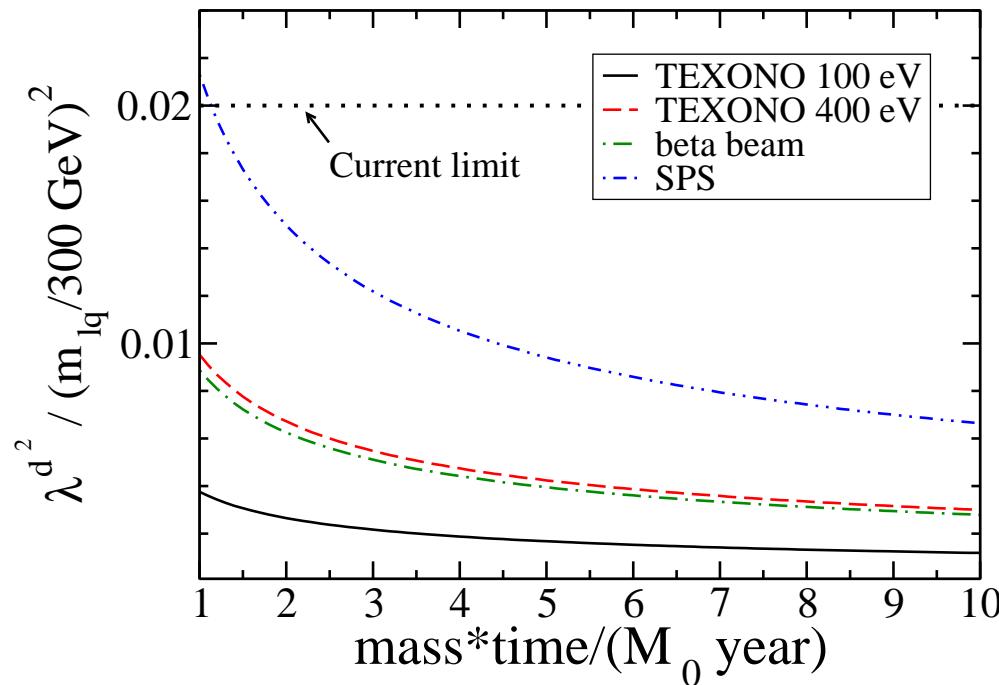
$$\begin{aligned}\varepsilon_{ee}^{uL} &= \varepsilon_{ee}^{dL} = -\varepsilon_{ee}^{uR} = -4\gamma \sin^2 \theta_W \rho_{\nu N}^{NC} \left(\frac{c_\beta}{\sqrt{24}} - \frac{s_\beta}{3} \sqrt{\frac{5}{8}} \right) \left(\frac{3c_\beta}{2\sqrt{24}} + \frac{s_\beta}{6} \sqrt{\frac{5}{8}} \right) \\ \varepsilon_{ee}^{dR} &= -8\gamma \sin^2 \theta_W \rho_{\nu N}^{NC} \left(\frac{3c_\beta}{2\sqrt{24}} + \frac{s_\beta}{6} \sqrt{\frac{5}{8}} \right)^2, \quad \gamma = (M_Z/M_{Z'})^2\end{aligned}$$



Leptoquark

$$\varepsilon^{uV} = \frac{\lambda_u^2}{m_{lq}^2} \frac{\sqrt{2}}{4G_F}, \quad \varepsilon^{dV} = \frac{\lambda_d^2}{m_{lq}^2} \frac{\sqrt{2}}{4G_F}$$

See Davidson et al'94



In some models leptoquark can be light

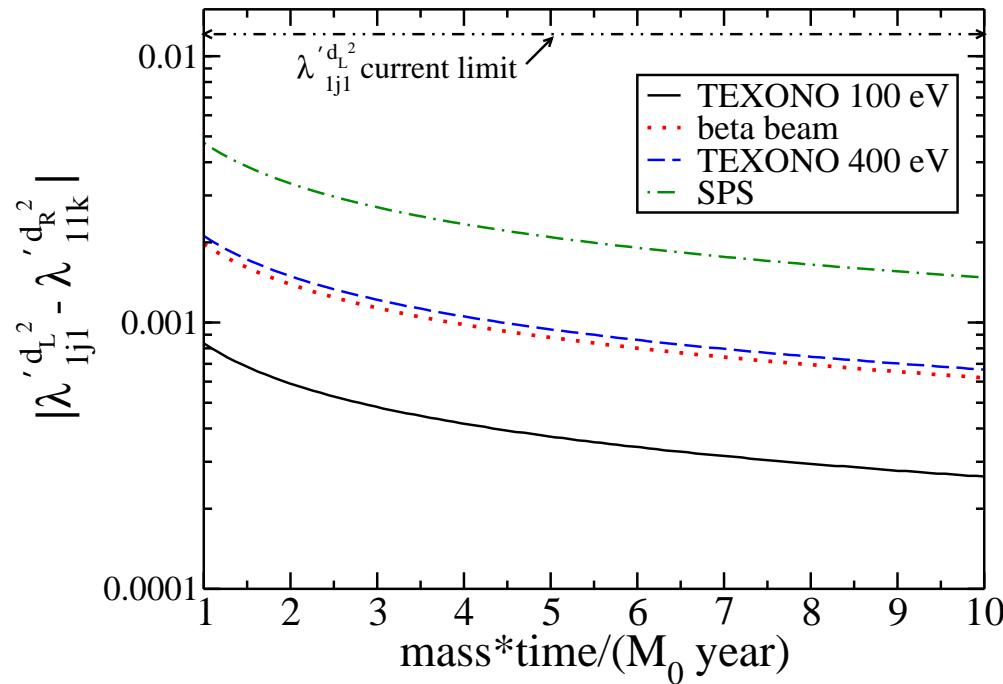
See Dorsner et al'05

SUSY with broken R-parity

R-parity violating MSSM (imposing baryon number conservation) with a superpotential that contains the following L - violating terms:

See Barger et al'89

$$\lambda_{ijk} L_L^i L_L^j \bar{E}_R^k, \quad \lambda'_{ijk} L_L^i Q_L^j \bar{D}_R^k$$



Conclusions

- Non-standard interactions in the neutrino sector are predicted in models beyond the SM and can be rather large
- NSI can play significant role in astrophysical environments
- Coherent neutrino scattering off nuclei gives a new sensitive probe to ν -quark vector couplings
- Degeneracy in determination of NSI parameters can be successfully resolved by taking two different targets
- Estimated sensitivity to NSI couplings is comparable and better than those expected from future ν -Factory

neutrino-nuclei scattering

| | $^{76}\text{Ge} + ^{28}\text{Si}$ $T_{th}=400\text{eV}$ | $^{76}\text{Ge} + ^{28}\text{Si}$ $T_{th}=100\text{eV}$ |
|--------------------------|---|---|
| ϵ_{ee}^{dV} | $ \epsilon_{ee}^{dV} < 0.036$ | $ \epsilon_{ee}^{dV} < 0.018$ |
| ϵ_{ee}^{uV} | $ \epsilon_{ee}^{uV} < 0.038$ | $ \epsilon_{ee}^{uV} < 0.019$ |
| $\epsilon_{\tau e}^{dV}$ | $ \epsilon_{\tau e}^{dV} < 0.48$ | $ \epsilon_{\tau e}^{dV} < 0.34$ |
| $\epsilon_{\tau e}^{uV}$ | $ \epsilon_{\tau e}^{uV} < 0.50$ | $ \epsilon_{\tau e}^{uV} < 0.37$ |
| ϵ_{ee}^{dV} | $-0.002 < \epsilon_{ee}^{dV} < 0.034$ | $-0.0009 < \epsilon_{ee}^{dV} < 0.016$ |
| $\epsilon_{\tau e}^{dV}$ | $ \epsilon_{\tau e}^{dV} < 0.1$ | $ \epsilon_{\tau e}^{dV} < 0.074$ |

NSI with d-quark only

$\varepsilon_{\tau e}^{dV}$ versus ε_{ee}^{dV}

